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## Crystal Structure

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## $\alpha$-Tris(2,4-pentanedionato- $\kappa^{2} O, O^{\prime}$ )chromium(III) at 290 and 110 K : a new $\delta$ phase at 110 K

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The crystal structure of the title compound, $\left[\mathrm{Cr}\left(\mathrm{C}_{5} \mathrm{H}_{7} \mathrm{O}_{2}\right)_{3}\right]$, has been determined at 290 and 110 K to provide information on thermal vibrations and disorder. The $\alpha$ polymorph at room temperature has been reported [Morosin (1965). Acta Cryst. 19, 131-137]. The reinvestigation of this structure, presented here, indicates the presence of weak uninterpretable supercell reflections together with disorder streaks. The discussed structure can thus be considered as an average structure. After cooling to 110 K , a new $\delta$ polymorph was found, which is a superstructure of the $\alpha$ polymorph. The space group remains $P 2_{1} / c$ and the phase transition can therefore be considered as klassengleich. The unit-cell volume increases by a factor of six, resulting in six independent molecules in the asymmetric unit.

## Comment

Metal atoms, especially transition metals, are present in the active centers of almost one-third of all known enzymes and play an important role in catalytic processes. $\mathrm{Cr}^{\text {III }}$ is an essential trace element (Goldhaber, 2003) with the important function of maintaining the normal glucose metabolism in the human body, in particular at advanced age or in the case of diabetes. Trivalent metal acetylacetonate complexes [ $M^{\mathrm{III}}(\mathrm{acac})_{3}$; acac is acetylacetonate or 2,4-pentanedionate] belong to the $\beta$-diketonates and are a particularly accessible species for studying structure, bonding and ligand coordination in organometallic systems. Owing to their high volatility, $\beta$-diketonates of metals are used as precursors in the process of chemical deposition of coatings from the vaporous phase, which is known to be a promising technique for the production of, for example, superconducting materials and nano-sized coatings (Naumov et al., 2006).

It is well known that these complexes show polymorphism; an overview is given by Sabolović et al. (2004). Phase transitions of these compounds are observed at low temperatures; for $\mathrm{Mn}(\mathrm{acac})_{3}$, see Geremia \& Demitri (2005), and for $\mathrm{Al}(\mathrm{acac})_{3}$, see von Chrzanowski et al. (2007a). $M^{\mathrm{III}}(\mathrm{acac})_{3}$ can
be grouped into a polymorphous series of three. By an old nomenclature of Astbury \& Morgan (1926), the $\alpha$ polymorph crystallizes in the monoclinic crystal system with the space group $P 2_{1} / c$. To our knowledge, all previous studies of $\mathrm{Cr}(\mathrm{acac})_{3}$, (I), have been carried out at room temperature, with the compound crystallizing as its $\alpha$ polymorph (e.g. Morosin, 1965). In the known studies of the $\alpha$ polymorphs, the C atoms show anomalous displacement parameters and disorder, especially in one of the three ligands (e.g. Hon \& Pfluger, 1973). In the course of our ongoing studies of trivalent metal acetylacetonates and in order to investigate this behavior, we report here the redetermination of $\alpha-\mathrm{Cr}(\mathrm{acac})_{3}$ at 290 K , ( $\mathrm{I} a$ ), and a new determination at 110 K , (I $b$ ). At 110 K , we observed a new $\delta$ phase.


As mentioned previously, $M^{\mathrm{III}}(\mathrm{acac})_{3}$ complexes show anomalous displacement parameters and disorder, especially in one of the three acac ligands, labeled $\mathrm{C} 11-\mathrm{C} 15$. We reported this observation in our previous temperature-dependent studies of $\alpha-\mathrm{Al}(\mathrm{acac})_{3}$ (von Chrzanowski et al., 2007a) and $\alpha-\mathrm{Co}(\mathrm{acac})_{3}$ (von Chrzanowski et al., 2007b). This can also be seen in the room-temperature structure of (I) reported by Morosin (1965), ( $\mathrm{I} m$ ), as well as in the present study of ( $\mathrm{I} a$ ) at 290 K (Figs. $1 a$ and $1 b$ ). Rigid-body analyses for both structures (Schomaker \& Trueblood, 1998) result in relatively high agreement factors of $R=0.204$ for ( $\mathrm{I} m$ ) and $R=0.131$ for ( $\mathrm{I} a$ ) $\left[R=\left\{\sum\left[w\left(U_{\text {obs }}-U_{\text {calc }}\right)^{2} / \sum\left(w U_{\text {obs }}\right)^{2}\right\}^{1 / 2}\right]\right.$. As a result of different weighting, the agreement factor for ( $\mathrm{I} m$ ) is significantly higher; for ( $\mathrm{I} a$ ), the weights were derived from the standard uncertainties of the anisotropic displacement parameters, while for ( $\mathrm{I} m$ ), unit weights were used because no standard uncertainties are given in the paper. Difference plots (Hummel et al., 1990) between the observed displacement parameters and the rigid-body models indicate large internal motions (Figs. $1 c$ and $1 d$ ).

This anomalous behavior is also manifested in the CCD detector images of ( $\mathrm{I} a$ ) by the presence of diffuse disorder streaks between the reflections (Fig. 2), on which some intensity maxima are present. These maxima were not interpreted because of their weakness. The presence of large isolated diffuse spots/areas was also reported in (I $m$ ) and there explained by thermal motion. On the basis of our data, we assume that ( $\mathrm{I} a$ ) actually has a larger unit cell, which

(a)

(b)

(c)

(d)

Figure 1
Left: displacement ellipsoid plots and atomic numbering schemes of (I), showing (a) the literature structure of (Im) (Morosin, 1965) at room temperature with the original atomic numbering scheme and $(b)(\mathrm{I} a)$ at 290 K . Ellipsoids are drawn at the $50 \%$ probability level and H atoms have been omitted for clarity. Right: peanut plots (Hummel et al., 1990) showing the difference between the measured displacement parameters and the parameters obtained by rigid-body analyses using the program THMA11 (Schomaker \& Trueblood, 1998). A scale factor of 3.08 was used for the r.m.s. surfaces. [In the electronic version of the paper, blue spheres (darker here) indicate positive differences and purple spheres (lighter here) negative differences.] (c) The literature structure of Morosin (1965) at room temperature and $(d)(\mathrm{I} a)$ at 290 K . The plots are drawn in the same orientation as the displacement ellipsoid plots and have the same atomic numbering schemes.


Figure 2
A CCD detector image of $(\mathrm{I} a)$ at 290 K , showing diffuse disorder streaks between reflections. Predicted reflections of the 1706.01 (18) $\AA^{3}$ unit cell are circled.
corresponds to the low-temperature $\delta$ polymorph (see below), and by ignoring the weak supercell reflections an average structure was obtained.

At 110 K , the diffraction pattern shows supercell reflections, which are now interpretable, and the structure of a new $\delta$ polymorph (I $b$ ) could be determined. Because this supercell is probably already present in ( $\mathrm{I} a)$, an exact temperature range for the phase transition cannot be determined. The transformation matrix from the high-temperature $\alpha$ polymorph to the new low-temperature $\delta$ polymorph is (200/010/003). The determinant of this matrix is six and therefore the cell volume increases by a factor of six. Interestingly, the application of this transformation matrix to the coordinates of ( $\mathrm{I} a$ ) at the chosen origin to generate the coordinates of ( $\mathrm{I} b$ ) leads to a false minimum. Solution attempts of (Ib) from scratch with the programs SIR97 (Altomare et al., 1999), SIR2000 (Burla et al., 2000) or DIRDIF99 (Beurskens et al., 1999) using standard parameters led to the same false minimum. In the false minimum, the structure can be refined to acceptable agreement factors, but the displacement parameters are unexpectedly large. The correct minimum was found by structure solution with the program SHELXS97 (Sheldrick, 1997) using direct methods and corresponds to an origin shift of $\frac{1}{4}$ in the direction of the crystallographic $a$ axis (Fig. 3). Refinement with the correct origin results in significantly improved agreement factors and, more importantly, completely normal
displacement parameters. With respect to the symmetry elements, the origin shift of $\frac{1}{4}$ means that exact inversion centers and $2_{1}$ screw axes in ( $\mathrm{I} a$ ) become pseudosymmetry elements in (Ib). In (Ib), the space group remains $P 2_{1} / c$, and the phase transition can thus be classified as klassengleich. A similar phase transition was observed in $\alpha-\mathrm{Al}(\mathrm{acac})_{3}$ (von Chrzanowski et al., 2007a). There, the phase transition occurs between 150 and 110 K , and the transformation matrix from the high-temperature $\alpha$ polymorph to the low-temperature $\delta$ polymorph is $(10 \overline{1} / 0 \overline{1} 0 / \overline{2} 0 \overline{1})$; the determinant of the matrix is three and the volume consequently increases by a factor of three.

The asymmetric unit of ( $\mathrm{I} b)$ at 110 K consists of six independent molecules (Fig. 4). All molecules have approximately noncrystallographic $D_{3}$ symmetry, with r.m.s. deviations from ideal symmetry of between $0.128 \AA$ for molecule 4 and 0.199 Å for molecule 5 (Pilati \& Forni, 1998). All six molecules have essentially the same geometry, as can be seen in a quaternion fit (Fig. 5a). This quaternion fit (Mackay, 1984) considers only the molecular structures but does not take the crystal packing into account. The packing effects can be seen by the application of the transformation matrix to the $\alpha$ polymorph on the atomic coordinates of the $\delta$ polymorph. The result of this operation can be seen in Fig. 5(b). Two of the acac ligands have only small deviations after this averaging,


Figure 3
The crystal packing of ( $\mathrm{I} b$ ) at 110 K , viewed along the crystallographic $b$ axis. The twofold screw axes are shown with black symbols and inversion centers are shown with unshaded circles. The unit cell of the hightemperature $\alpha$ phase is drawn with dashed lines; symbols for the screw axes and the inversion centers are shown shaded. Molecules of ( $\mathrm{I} a$ ) are not displayed. [In the electronic version of the paper, for the molecules of $(\mathrm{I} b)$, orange represents molecule $1(X=1)$, yellow molecule $2(X=2)$, blue molecule $3(X=3)$, green molecule $4(X=4)$, red molecule $5(X=5)$ and black molecule $6(X=6)$.]
while the third ligand ( $\mathrm{C} 11 X-\mathrm{C} 15 X$ ) is severely affected by the packing. The latter ligand corresponds to the ligand with large displacement parameters (C11-C15) in the $\alpha$ polymorph, and therefore the phase transition can be considered as a disorderorder phase transition.

A possible explanation for this behavior is the presence of different $\mathrm{C}-\mathrm{H} \cdots \mathrm{O}$ interactions (Table 3). Every ligand is a donor of one intermolecular $\mathrm{C}-\mathrm{H} \cdots \mathrm{O}$ interaction with a range of $\mathrm{C}-\mathrm{H} \cdots \mathrm{O}$ angles between 163 and $179^{\circ}$. The average $\mathrm{C} \cdots \mathrm{O}$ distance for the $\mathrm{C} 11 X-\mathrm{C} 15 X$ ligand $(X=1-6$ for molecules 1-6) is 3.57 (2) $\AA$ (for $\mathrm{C} 15 X-\mathrm{H} \cdots \mathrm{O} 4 Y$ ); this is the longest contact and therefore the weakest $\mathrm{C}-\mathrm{H} \cdots \mathrm{O}$ interaction, resulting in the largest deviation for $\mathrm{C} 11 X-\mathrm{C} 15 X$ (Table 4). The ligands appear to be split into two different groups; molecules 1, 2 and 4 define group 1, and molecules 3, 5 and 6 define group 2 (Table 5).

The ADDSYM routine of the program PLATON (Spek, 2003) indicates pseudo-translational symmetry for (Ib) at 110 K . Such a situation has also been observed in the $\delta$ phase of $\mathrm{Al}(\mathrm{acac})_{3}$ (von Chrzanowski et al., 2007a), where two groups of reflections were present, viz. strong subcell reflections and weak supercell reflections. In (Ib), three groups of reflections are present, viz. strong subcell reflections, and weak and very weak supercell reflections (Table 6). Nevertheless the weak supercell reflections are clearly present and prevent a transformation to the subcell.

A pseudo-translational symmetry test (SIR97; Altomare et al., 1999) results in a value of $74 \%$ for the mean fractional scattering power of the electron density for reflections with $h=$ $2 n$ and $l=3 n\left(\left\langle E^{2}\right\rangle=4.734\right)$, which is based on normalized structure factors from measured data (Cascarano et al., 1985). The three highest vectors of nonzero length in the Patterson map, as calculated with SHELXS97 (Sheldrick, 1997), end at $(0.4994,0.0000,0.3332),(0.9996,0.0000,0.3332)$ and $\left(\frac{1}{2}, 0,0\right)$, again confirming the pseudo-translational symmetry of $\frac{1}{2}$ in the $a$ and $\frac{1}{3}$ in the $c$ direction.

Despite the pseudo-translational symmetry, a full matrix least-squares refinement with SHELXL97 (Sheldrick, 1997) can be performed with default refinement parameters and without restraints or constraints. No correlation matrix elements were larger than 0.5 . The weighting scheme for the refinement was optimized by SHELXL97 based on all 18651 unique reflections and results in a goodness-of-fit of 1.064. The corresponding goodness-of-fit for the 3155 strong subcell reflections without rerefinement and with the same weighting scheme is 1.748 , and those for the 6296 weak and 9320 very weak supercell reflections are 1.038 and 1.029 , respectively. Obviously the weighted $\sigma$ values of the subcell reflections are underestimated.

A distribution diagram of all $\mathrm{Cr}-\mathrm{O}$ bond distances shows a single maximum, and therefore the averaging of these distances is valid. A normal probability plot (Abrahams \& Keve, 1971) of the low-temperature $\delta$ polymorph, in which all bond distances of each molecule are treated independently, results in a slope of 1.965 . The observed outliers belong to the $\mathrm{Cr}-\mathrm{O}$ distances. The range of $\mathrm{Cr}-\mathrm{O}$ distances is wide; they range from 1.9460 (16) to 1.9750 (17) $\AA$ (Table 2). There is no
chemical evidence for this large variance. This effect can be attributed to the internal motion of the molecules and the high standard uncertainties of the $\mathrm{Cr}-\mathrm{O}$ bond distances. $\gamma-\mathrm{Al}(\mathrm{acac})_{3}$ (von Chrzanowski et al., 2006) and $\delta-\mathrm{Al}(\mathrm{acac})_{3}$ (von Chrzanowski et al., 2007a) show similar variations of the corresponding $\mathrm{Al}-\mathrm{O}$ bond distances, which range from
1.8790 (12) to $1.8909(13) \AA$ for $\gamma-\mathrm{Al}(\mathrm{acac})_{3}$ and from 1.8704 (9) to 1.8910 (10) $\AA$ for $\delta-\mathrm{Al}(\mathrm{acac})_{3}$. The average values of $\mathrm{Cr}-\mathrm{O}$ distances of the six independent molecules can be grouped into long $\mathrm{Cr}-\mathrm{O}$ distances for molecules 2,3 and 4 [1.9683 $(\mathrm{Cr} X-\mathrm{O} 4 X)$ to $1.9673 \AA(\mathrm{Cr} X-\mathrm{O} 2 X)]$ and short distances for molecules 1,5 and $6[1.9536(\mathrm{Cr} X-\mathrm{O} 6 X)$ to

(a)

(d)

(b)

(e)

(c)

(f)

Figure 4
Displacement ellipsoid plots and atomic numbering schemes of the six independent molecules of ( $\mathrm{I} b$ ) at 110 K . All molecules are shown independently using the same orientation with the same atomic numbering scheme and do not represent the crystal packing. Ellipsoids are drawn at the $50 \%$ probability level. H atoms have been omitted for clarity. The views show $(a)$ molecule $1(X=1),(b)$ molecule $2(X=2),(c)$ molecule $3(X=3)$, $(d)$ molecule $4(X=$ 4), (e) molecule $5(X=5)$ and $(f)$ molecule $6(X=6)$.

(a)

(b)

Figure 5
(a) Quaternion fit (Mackay, 1984) of the six independent molecules of ( $\mathrm{I} b$ ) at 110 K . (b) Transformation of the six independent molecules of ( $\mathrm{I} b$ ) at 110 K into an average structure using the transformation matrix between the $\alpha$ and $\delta$ polymorphs. (The color scheme for the electronic version is the same as in Fig. 3.)
$1.9589 \AA(\mathrm{Cr} X-\mathrm{O} 1 X)]$. Each long averaged $\mathrm{Cr}-\mathrm{O}$ distance corresponds to a short distance on the opposite side.

This variation is also present in the high-temperature $\alpha$ polymorph ( $\mathrm{I} a$ ), with $\mathrm{Cr}-\mathrm{O}$ bond distances ranging between 1.9413 (18) and 1.9645 (17) $\AA$ (Table 1). The corresponding $M-\mathrm{O}$ distances in $\alpha-\mathrm{Al}(\mathrm{acac})_{3}$ and $\alpha-\mathrm{Co}(\mathrm{acac})_{3}$ show a similar variation (von Chrzanowski et al., 2007a,b). Because octahedral $\mathrm{Cr}^{\mathrm{III}}$ with electron configuration $d^{3}$ does not exhibit Jahn-Teller distortions (Wiberg, 1985), this variation can again only be explained by internal thermal motion. The thermal motion also contributes to a shortening of the $\mathrm{Cr}-\mathrm{O}$ distances of the crystal structure determinations compared with $1.9773 \AA$ obtained for $\mathrm{Cr}(\mathrm{acac})_{3}$ from high-level density functional theory (DFT) (B3LYP/6-31G* for C, H, O and triple- $\zeta$ for Cr ) calculations (Diaz-Acosta et al., 2001).

## Experimental

Violet crystals were obtained by slow evaporation of a solution of the commercially available material (Aldrich) in ethyl acetate at room temperature.

## Determination (la)

## Crystal data

| $\left[\mathrm{Cr}\left(\mathrm{C}_{5} \mathrm{H}_{7} \mathrm{O}_{2}\right)_{3}\right]$ | $V=1706.01(18) \AA^{3}$ |
| :--- | :--- |
| $M_{r}=349.32$ | $Z=4$ |
| Monoclinic, $P 2_{1} / c$ | Mo $\alpha \alpha$ radiation |
| $a=13.9990(8) \AA$ | $\mu=0.69 \mathrm{~mm}^{-1}$ |
| $b=7.5441(4) \AA$ | $T=290(2) \mathrm{K}$ |
| $c=16.3590(12) \AA$ | $0.30 \times 0.27 \times 0.12 \mathrm{~mm}$ |
| $\beta=99.031(2)^{\circ}$ |  |
| Data collection |  |
| Nonius KappaCCD diffractometer | 45735 measured reflections |
| Absortion correction: multi-scan | 3924 independent reflections |
| $($ SADABS; Sheldrick, 2002 $)$ | 3135 reflections with $I>2 \sigma(I)$ |
| $T_{\text {min }}=0.68, T_{\text {max }}=0.92$ | $R_{\text {int }}=0.030$ |

## Refinement

$R\left[F^{2}>2 \sigma\left(F^{2}\right)\right]=0.042$
$w R\left(F^{2}\right)=0.127$
$S=1.04$
3924 reflections

## Determination (lb)

## Crystal data

$\left[\mathrm{Cr}_{\left.\left(\mathrm{C}_{5} \mathrm{H}_{7} \mathrm{O}_{2}\right)_{3}\right]}\right.$
$V=9877.09(17) \AA^{3}$
$M_{r}=349.32$
Monoclinic, $P 2_{1} / c$
$a=27.5823$ (2) A
$b=7.4656$ (1) $\AA$
$c=48.6012$ (4) $\AA$
$\beta=99.2731$ (3) ${ }^{\circ}$

## Data collection

Nonius KappaCCD diffractometer Absorption correction: multi-scan (SORTAV; Blessing, 1987)
$T_{\text {min }}=0.82, T_{\text {max }}=0.94$

## Refinement

$R\left[F^{2}>2 \sigma\left(F^{2}\right)\right]=0.040$
$w R\left(F^{2}\right)=0.139$
$S=1.06$
18651 reflections
$Z=24$
Mo $K \alpha$ radiation
$\mu=0.72 \mathrm{~mm}^{-1}$
$T=110$ (2) K
$0.39 \times 0.29 \times 0.09 \mathrm{~mm}$

101596 measured reflections 18651 independent reflections 9662 reflections with $I>2 \sigma(I)$ $R_{\text {int }}=0.052$

1225 parameters
H -atom parameters constrained
$\Delta \rho_{\text {max }}=0.35$ e $\AA^{-3}$
$\Delta \rho_{\text {min }}=-0.60 \mathrm{e}^{-3}$

Table 3
Hydrogen-bond geometry ( $\mathrm{A},{ }^{\circ}$ ) for (Ib).

| $D-\mathrm{H} \cdots A$ | $D-\mathrm{H}$ | $\mathrm{H} \cdots A$ | $D \cdots A$ | $D-\mathrm{H} \cdots A$ |
| :---: | :---: | :---: | :---: | :---: |
| C51-H51C $\cdots$ O52 | 0.98 | 2.49 | 3.465 (3) | 172 |
| C52-H52B $\cdots$ O53 | 0.98 | 2.45 | 3.423 (3) | 173 |
| $\mathrm{C} 53-\mathrm{H} 53 \mathrm{C} \cdots \mathrm{O}^{1}{ }^{\text {i }}$ | 0.98 | 2.52 | 3.470 (3) | 163 |
| C54-H54B $\cdots \mathrm{O} 5$ | 0.98 | 2.46 | 3.438 (3) | 173 |
| C55-H55C $\cdots \mathrm{O} 56$ | 0.98 | 2.49 | 3.458 (3) | 169 |
| C56-H56B $\cdots$ O54 ${ }^{\text {ii }}$ | 0.98 | 2.52 | 3.468 (3) | 163 |
| $\mathrm{C} 101-\mathrm{H} 10 \mathrm{C} \cdots \mathrm{O} 11^{\text {iii }}$ | 0.98 | 2.60 | 3.584 (3) | 179 |
| $\mathrm{C} 102-\mathrm{H} 10 \mathrm{D} \cdots \mathrm{O} 12^{\text {iv }}$ | 0.98 | 2.60 | 3.582 (3) | 177 |
| C103-H10I $\cdots$ O13 ${ }^{\text {iii }}$ | 0.98 | 2.53 | 3.509 (3) | 176 |
| C104-H10J . . O14 ${ }^{\text {iv }}$ | 0.98 | 2.61 | 3.591 (3) | 178 |
| $\mathrm{C} 105-\mathrm{H} 10 \mathrm{O} \cdots \mathrm{O} 15^{\text {iii }}$ | 0.98 | 2.56 | 3.533 (3) | 173 |
| C106-H10P ...O16 ${ }^{\text {iv }}$ | 0.98 | 2.52 | 3.493 (3) | 175 |
| $\mathrm{C} 151-\mathrm{H} 15 \mathrm{~B} \cdots \mathrm{O} 46^{\text {iii }}$ | 0.98 | 2.61 | 3.580 (3) | 172 |
| C152-H15E . . ${ }^{\text {4 }} 4$ | 0.98 | 2.61 | 3.587 (3) | 172 |
| C153-H15I . . O44 ${ }^{\text {iii }}$ | 0.98 | 2.59 | 3.550 (3) | 165 |
| C154-H15K...O43 ${ }^{\text {iv }}$ | 0.98 | 2.62 | 3.591 (3) | 171 |
| C155-H15R $\ldots$ O 42 | 0.98 | 2.56 | 3.531 (3) | 169 |
| C156-H15T . . O44 ${ }^{\text {iv }}$ | 0.98 | 2.60 | 3.562 (3) | 166 |

Symmetry codes: (i) $x,-y+\frac{1}{2}, z+\frac{1}{2}$; (ii) $x,-y+\frac{3}{2}, z-\frac{1}{2}$; (iii) $x, y-1, z$; (iv) $x, y+1, z$.

The X-ray intensities of (I $a$ ) were obtained with two different exposure times and rotation angles of $1^{\circ}$ at $290 \mathrm{~K} .364 \varphi$ and $455 \omega$ scans were measured with an exposure time of 60 s per frame, and 364 $\varphi$ scans with an exposure time of 12 s per frame. The X-ray intensities of ( $\mathrm{I} b$ ) were measured with a different crystal at 110 K and were obtained with two different data sets. For data set $1,1383 \varphi$ scans and $1114 \omega$ scans were measured with an exposure time of 10 s per frame and a rotation angle of $0.5^{\circ}$ (overflow scans).

Table 2
Selected bond lengths ( $(\AA)$ for the molecules of ( $\mathrm{I} b$ ).

|  | $X=1$ | $X=2$ | $X=3$ | $X=4$ | $X=5$ | $X=6$ | $d_{\max }-d_{\min }$ |
| :--- | :--- | :--- | :--- | :--- | :--- | :--- | :--- |
| Cr1-OX1 | $1.9572(17)$ | $1.9580(17)$ | $1.9690(16)$ | $1.9703(16)$ | $1.9651(18)$ | $1.9437(17)$ | 0.0266 |
| Cr2-O 22 | $1.9541(16)$ | $1.9705(17)$ | $1.9662(16)$ | $1.9692(16)$ | $1.9569(17)$ | $1.9481(16)$ | 0.0224 |
| Cr3-OX3 | $1.9614(17)$ | $1.9645(18)$ | $1.9676(16)$ | $1.9700(17)$ | $1.9564(17)$ | $1.9623(16)$ | 0.0136 |
| Cr4-OX4 | $1.9640(17)$ | $1.9685(18)$ | $1.9696(16)$ | $1.9704(17)$ | $1.9663(17)$ | $1.9460(16)$ | 0.0244 |
| Cr5-OX5 | $1.9602(17)$ | $1.9672(17)$ | $1.9673(17)$ | $1.9670(17)$ | $1.9561(18)$ | $1.9567(17)$ | 0.0166 |
| Cr6-OX6 | $1.9566(17)$ | $1.9750(17)$ | $1.9682(16)$ | $1.9629(17)$ | $1.9501(17)$ | $1.9645(16)$ | 0.0249 |
| Average | 1.9589 | 1.9673 | 1.9680 | 1.9683 | 1.9585 | 1.9536 |  |

Note: $X=1-6$ for molecules 1-6.

Table 4
Averaged $\mathrm{C}-\mathrm{H} \cdots \mathrm{O}$ interactions ( $\AA$ ) of ligands $\mathrm{C} 1 X-\mathrm{C} 5 X, \mathrm{C} 5 X-\mathrm{C} 10 X$ and C11X-C15X for (Ib).

| $D-\mathrm{H} \cdots A$ | $\mathrm{H} \cdots A$ | $D \cdots A$ |
| :--- | :--- | :--- |
| C5 $X-\mathrm{H} \cdots \mathrm{O} Y$ | 2.49 | $3.45(2)$ |
| C10 $X-\mathrm{H} \cdots \mathrm{O} 1 Y$ | 2.57 | $3.55(4)$ |
| C15X-H $\cdots$ O4Y | 2.60 | $3.57(2)$ |

Notes: calculation of the average values used [ $\left.\sum \mathrm{C} X-\mathrm{H} X Y \cdots \mathrm{O} X\right) / 6$ ], with $X=1-6$ for molecules 1-6. $Y$ takes labels and symmetry codes from Table 3.

Table 5
Averaged $\mathrm{C}-\mathrm{H} \cdots \mathrm{O}$ interactions ( A ) for molecules 1-6 of (Ib).

| Group | Molecule | $\mathrm{H} \cdots A$ | $D \cdots A$ |
| :--- | :--- | :--- | :--- |
| 1 | 1 | 2.57 | $3.54(7)$ |
| 1 | 2 | 2.56 | $3.53(9)$ |
| 1 | 4 | 2.56 | $3.54(9)$ |
| 2 | 3 | 2.55 | $3.51(4)$ |
| 2 | 5 | 2.54 | $3.51(4)$ |
| 2 | 6 | 2.54 | $3.51(5)$ |

Note: calculation of the average values used $[(\mathrm{C} 5 X-\mathrm{H} 5 X \cdots \mathrm{O} 5 X+\mathrm{C} 10 X-$ $\mathrm{H} 10 X \cdots \mathrm{O} 1 X+\mathrm{C} 15 X-\mathrm{H} 15 X \cdots \mathrm{O} 4 X) / 3]$, with $X=1-6$ for molecules 1-6.

Table 6
Intensity statistics in (Ib).

|  | Reflections | $\langle I\rangle$ | $\langle\sigma\rangle$ | $\langle I\rangle /\langle\sigma\rangle$ | $\langle I / \sigma\rangle$ |
| :--- | :--- | :--- | :--- | :--- | :--- |
| $A$ | 3115 | 23111.90 | 467.53 | 49.43 | 32.17 |
| $B$ | 6216 | 1652.16 | 95.29 | 17.34 | 13.73 |
| $C$ | 9320 | 138.88 | 59.67 | 2.33 | 2.31 |

Notes: $A: h=2 n$ and $l=3 m$ ( $m$ and $n$ are integers); $B: h \neq 2 n$ and $l \neq 3 m$ ( $m$ and $n$ are integers); $C$ : neither $A$ nor $B$.

The data were integrated with EVAL14 (Duisenberg et al., 2003). Only reflections with $I>10 \sigma(I)$ were used. For data set 2 , the same experiment was repeated with an exposure time of 50 s per frame. DENZO and SCALEPACK (Otwinowski \& Minor, 1997) were used for the integration of the intensities of all data. The two data sets were scaled and corrected for absorption with SORTAV (Blessing, 1987). All H atoms were introduced in geometrically idealized positions ( $\mathrm{C}-\mathrm{H}=0.93-0.98 \AA$ ), refined with a riding model and subsequently confirmed in contoured difference Fourier maps. Their isotropic displacement parameters were constrained $\left[U_{\text {iso }}(\mathrm{H})=1.2 U_{\text {eq }}(\mathrm{C})\right.$ for H atoms of the central $\mathrm{C}-\mathrm{H}$ groups and $U_{\text {iso }}(\mathrm{H})=1.5 U_{\text {eq }}(\mathrm{C})$ for methyl H atoms].

For (I $a$ ), data collection: COLLECT (Nonius, 1999); cell refinement: PEAKREF (Schreurs, 2005); data reduction: EVAL14 (Duisenberg et al., 2003) and SADABS (Sheldrick, 2002); program(s) used to solve structure: coordinates taken from the isostructural $\alpha-\mathrm{Co}(\mathrm{acac})_{3}$ compound (von Chrzanowski et al., 2007b). For (Ib), data collection: COLLECT (Nonius, 1999); cell refinement: HKL2000 (Otwinowski \& Minor, 1997); data reduction: EVAL14 (set 1) (Duisenberg et al., 2003), HKL-2000 (set 2) and SORTAV (Blessing, 1987); program(s) used to solve structure: SHELXS97 (Sheldrick, 1997). For both compounds, program(s) used to refine structure:

SHELXL97 (Sheldrick, 1997); molecular graphics: PLATON (Spek, 2003); software used to prepare material for publication: manual editing of SHELXL97 output.

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Supplementary data for this paper are available from the IUCr electronic archives (Reference: BG3044). Services for accessing these data are described at the back of the journal.

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